Modeling and Simulations of Light Emission and Propagation in Open Nanophotonic Systems



Jakob Rosenkrantz de Lasson

Nanophotonics Theory & Signal Processing

www.nanophotonics.dk

Ph.D. Lecture, December 16 2015







I. Why





II. How





II. How

III. What



I. Why do we do our research?

















Vidensklumme: Verdens største maskine kører på lys Internettet vokser med 35-50 procent om året, og havde det ikke været for lyset, ville det være brudt sammen for længe siden.





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DTU Fotonik Department of Photonics Engineering



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DTU Fotonik Department of Photonics Engineering [D. A. B. Miller, ICTON 2009]



DTU Fotonik Department of Photonics Engineering [D. A. B. Miller, ICTON 2009]



[IBM]

[ETH Zürich]







NATURE|Vol 453|19 June 2008|doi:10.1038/nature07127

INSIGHT REVIEW

The quantum internet

H. J. Kimble¹

Quantum networks provide opportunities and challenges across a range of intellectual and technical frontiers, including quantum computation, communication and metrology. The realization of quantum networks composed of many nodes and channels requires new scientific capabilities for generating and characterizing quantum coherence and entanglement. Fundamental to this endeavour are quantum interconnects, which convert quantum states from one physical system to those of another in a reversible manner. Such quantum connectivity in networks can be achieved by the optical interactions of single photons and atoms, allowing the distribution of entanglement across the network and the teleportation of quantum states between nodes.

Forbes / Tech

DEC 13, 2012 @ 06:15 PM 8,905 VIEWS

Why IBM and Intel Are Chasing the \$100B Opportunity in Nanophotonics



D-Wave

Google's quantum computer is 100 million times faster than your laptop

But is it a true quantum computer?

DAVID NIELD 10 DEC 2015

DTU Fotonik Department of Photonics Engineering

desired



[C. Wilke, "The Serial Mentor"]















II. How do we do our research?

II. How do we do our research? – First half: The structures



THE ELECTROMAGNETIC SPECTRUM



THE ELECTROMAGNETIC SPECTRUM





[J. D. Joannopoulos et al., "Photonic Crystals – Molding the Flow of Light"]



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[Courtesy of Laboratory of Physics of Nanostructures, EPFL, Switzerland]

Lene Vestergaard Hau*†, S. E. Harris‡, Zachary Dutton*† & Cyrus H. Behroozi*§



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[T. F. Krauss, J. Phys. D 40, 2666 (2007)]



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Why do we need slow light?

Lene Vestergaard Hau*†, S. E. Harris‡, Zachary Dutton*† & Cyrus H. Behroozi*§

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Why do we need slow light?

"Because we can" Ü



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is so small that this process is not effective in bringing a spin system into thermal equilibrium with its surroundings. At 300°K, for $\nu = 10^{\circ}$ sec.⁻¹, $\mu = 1$ nuclear magneton, the corresponding relaxation time would be 5×10^{a1} seconds! However, for a system coupled to a resonant electrical circuit, the factor $8\pi^{a}/c^{3}$ no longer gives correctly the number of radiation oscillators per unit volume, in unit frequency range ν/Q associated with the circuit. The spontaneous emission probability is thereby increased, and the relaxation time reduced, by a factor $f = 3Q\lambda^{3}/a^{a2}V$, where V is the volume of the resonant. If a is a dimension characteristic of the circuit so that $V \sim a^{3}$, and if δ is the skin-depth at frequency $r, f \sim \lambda^{3}/a^{2}\delta$. For a non-resonant



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$$\Gamma = \frac{2\pi}{\hbar^2} \sum_{f} \left| \left\langle f \right| \hat{H}_{\mathrm{I}} \right| i \right\rangle \right|^2 \delta(\omega_i - \omega_f)$$
$$\hat{H}_{\mathrm{I}} = -\hat{\mathbf{p}} \cdot \hat{\mathbf{E}}$$

[L. Novotny and B. Hecht, "Principles of Nano-Optics" (2012)]

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II. How do we do our research? – First half: The structures

II. How do we do our research? – First half: The structures

II. *How* do we do our research? – First half: The structures

- Second half: The methods

Modeling and Simulations of Light Emission and Propagation in Open Nanophotonic Systems

Modeling and Simulations of Light Emission and Propagation in Open Nanophotonic Systems



["Stepping Stones to Mathematical Modeling", Indiana University]





 $\nabla \cdot \mathbf{D} = \rho$ $\nabla \cdot \mathbf{B} = 0$ $\nabla \times \mathbf{E} = -\partial_t \mathbf{B}$ $\nabla \times \mathbf{H} = \partial_t \mathbf{D} + \mathbf{J}$



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VIII. A Dynamical Theory of the Electromagnetic Field.

By J. CLERK MAXWELL, F.R.S.

[Phil. Trans. R. Soc. Lond. 155, 459-512 (1865)]





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DTU Indefra

Maxwells ligninger - sværere end kvantemekanik!

Af Jakob Rosenkrantz de Lasson 4. jun 2014 kl. 07:30

www.nanophotonics.dk







Von Felix Bloch in Leipzig.

Mit 2 Abbildungen. (Eingegangen am 10. August 1928.)





1952 Nobel Prize

Von Felix Bloch in Leipzig.

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III. What do we do in our research?

III. What do we do in our research? First half: Light control with slow light waveguides



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[M. Arcari et al., Phys. Rev. Lett. 113, 093603 (2014)]

See also: [V. S. C. Manga Rao and S. Hughes, Phys. Rev. B **75**, 205437 (2007)]
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[M. Arcari et al., Phys. Rev. Lett. 113, 093603 (2014)]







[E. Semenova et al., Appl. Phys. Lett. 99, 101106 (2011)]

Laboratory of Physics of Nanostructures (Eli Kapon), EPFL, Switzerland



[M. Calic et al., Phys. Rev. Lett. 106, 227402 (2011)]

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$$\frac{\rho_{\rm G}}{\rho_{\rm Bulk}} = \frac{3}{4\pi^2} \left(\frac{\lambda_0}{n}\right)^3 \frac{Q}{V_{\rm eff}} \eta \qquad \qquad Q = \frac{\omega a}{2c} n_{\rm G} \qquad \qquad V_{\rm eff} = \frac{\int_{\rm uc} \epsilon(\mathbf{r}) |\mathbf{E}(\mathbf{r})|^2 \, \mathrm{d}V}{\max\left[\epsilon(\mathbf{r})|\hat{\mathbf{n}} \cdot \mathbf{E}(\mathbf{r})|^2\right]}$$

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27/41







[Work in progress]





[Work in progress]

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[Work in progress]







The photonic band edge laser: A new approach to gain enhancement

Jonathan P. Dowling, Michael Scalora, Mark J. Bloemer, and Charles M. Bowden Weapont Sciences Directorate, AMSML-RD-WS, Research, Development, and Engineering Center, U. S. Army Missile Command, Reditore Arsenal, Alabama 35895-524

(Received 23 September 1993; accepted for publication 8 November 1993)

[J. P. Dowling et al., J. Appl. Phys. 75, 1896 (1994)]







[Y. Chen et al., Phys. Rev. A 92, 053839 (2015)]
 See also: [S. Ek et al., Nat. Commun. 5, 5039 (2014)]
 [J. Grgić et al., Phys. Rev. Lett. 108, 183903 (2012)]

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$$\partial_z \psi_+(z) = \frac{i\omega}{c} n_{gz} \chi_{\text{pert}}[\delta(z)\psi_+ + \kappa^*(z)e^{-i2k_z z}\psi_-],$$

$$\partial_z \psi_-(z) = -\frac{i\omega}{c} n_{gz} \chi_{\text{pert}}[\delta(z)\psi_- + \kappa(z)e^{i2k_z z}\psi_+],$$

[Y. Chen et al., Phys. Rev. A 92, 053839 (2015)]
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 [J. Grgić et al., Phys. Rev. Lett. 108, 183903 (2012)]







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Fundamental Limitations to Gain Enhancement in Periodic Media and Waveguides

Jure Grgić,¹ Johan Raunkjær Ott,¹ Fengwen Wang,² Ole Sigmund,² Antti-Pekka Jauho,³ Jesper Mørk,¹ and N. Asger Mortensen^{1,*}



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- Second half: Light control with cavities



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$A_{\nu} = (8\pi\nu^2/c^3)h\nu(8\pi^3\mu^2/3h^2)$ sec.⁻¹,

is so small that this process is not effective in bringing a spin system into thermal equilibrium with its surroundings. At 300°K, for $\nu = 10^{\circ}$ sec.⁻¹, $\mu = 1$ nuclear magneton, the corresponding relaxation time would be 5×10^{a1} seconds! However, for a system coupled to a resonant electrical circuit, the factor $8\pi r^{2}/c^{3}$ no longer gives correctly the number of radiation oscillators per unit volume, in unit frequency range , there being now one oscillator in the frequency range the associated with the circuit. The spontaneous emission probability is thereby increased, and the relaxation time reduced, by a factor $f = 3QN/4\pi^{2}V$, where *V* is the volume of the resonator. If *a* is a dimension characteristic of the circuit so that $V \sim a^{3}$, and if δ is the skin-depth at frequency $\nu, f \sim \lambda^{3}/a^{2}$. For a non-resonant circuit $f \sim \lambda^{3}/a^{3}$



c particles, of diameter 10^{-3} cm are mixed magnetic medium at room temperature, ission should establish thermal equilibrium order of minutes, for $\nu = 10^7$ sec.⁻¹.

[E. Purcell, Phys. Rev. 69, 681 (1946)]



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$$\frac{\Gamma(\mathbf{r};\omega)}{\Gamma_0} = \sum_{\mu} f_{\tilde{\omega}_{\mu}}(\mathbf{r};\omega)$$

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is so small that this process is not effective in bringing a spin system into thermal equilibrium with its surroundings. At 300°K, for $\nu = 10^{\circ}$ sec.⁻¹, $\mu = 1$ nuclear magneton, the corresponding relaxation time would be $5 \times 10^{\alpha_1}$ seconds! However, for a system coupled to a resonant electrical circuit, the factor $8\pi^{\mu}/c^{2}$ no longer gives correctly the number of radiation oscillators per unit volume, in unit frequency range , there being now one oscillator in the frequency range ν/Q associated with the circuit. The spontaneous emission probability is thereby increased, and the relaxation time reduced, by a factor $f=30\lambda^{3}/4\pi^{2}V$, where V is the volume of the resonator. If a is a dimension characteristic of the circuit so that $V \sim a^{3}$, and if δ is the skin-depth at frequency ν , $f \sim \lambda^{3}/a^{2}\delta$. For a non-resonant circuit $f \sim \lambda^{3}/a^{2}$, and for $a < \delta$ it can be shown that $f \sim \lambda^{3}/a^{2}$.



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$$\frac{\Gamma(\mathbf{r};\omega\simeq\tilde{\omega}_{\mu'})}{\Gamma_0}\simeq f_{\tilde{\omega}_{\mu'}}(\mathbf{r};\omega)\propto\frac{Q_{\mu'}}{V_{\mu'}}$$

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Quasi-normal modes







[P. T. Kristensen et al., Opt. Lett. 37, 1649 (2012)]





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[P. T. Kristensen et al., Opt. Lett. 39, 6359 (2014)]





[P. T. Kristensen et al., Opt. Lett. 39, 6359 (2014)]

[J. R. de Lasson et al., J. Opt. Soc. Am. A 31, 2142 (2014)]



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$$\Gamma(\mathbf{r};\omega) = \frac{\pi\omega}{\hbar\epsilon_0} |\mathbf{p}|^2 \rho(\mathbf{r};\omega)$$

[L. Novotny and B. Hecht,

"Principles of Nano-Optics" (2012)]



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$$\Gamma(\mathbf{r};\omega) = \frac{\pi\omega}{\hbar\epsilon_0} |\mathbf{p}|^2 \rho(\mathbf{r};\omega)$$
LDOS

$$\rho(\mathbf{r};\omega) = \frac{2\omega}{\pi c^2} \text{Im} \left[\mathbf{\hat{n}}_{\alpha} \cdot \mathbf{G}(\mathbf{r},\mathbf{r};\omega) \cdot \mathbf{\hat{n}}_{\alpha} \right]$$

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$$\Gamma(\mathbf{r};\omega) = \frac{\pi\omega}{\hbar\epsilon_0} |\mathbf{p}|^2 \frac{\rho(\mathbf{r};\omega)}{\text{LDOS}} \qquad \qquad \mathbf{G}(\mathbf{r},\mathbf{r}';\omega) = \frac{c^2}{2} \sum_{\mu} \frac{\mathbf{E}_{\mu}(\mathbf{r}) \otimes \mathbf{E}_{\mu}(\mathbf{r}')}{\tilde{\omega}_{\mu}(\tilde{\omega}_{\mu} - \omega)}$$

$$\rho(\mathbf{r};\omega) = \frac{2\omega}{\pi c^2} \operatorname{Im} \left[\mathbf{\hat{n}}_{\alpha} \cdot \mathbf{G}(\mathbf{r},\mathbf{r};\omega) \cdot \mathbf{\hat{n}}_{\alpha} \right]$$

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"Principles of Nano-Optics" (2012)]

See also: [C. Sauvan et al., Phys. Rev. Lett. **110**, 237401 (2013)] [R.-C. Ge et al., New J. Phys. **16**, 113048 (2014)]

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$$\rho(\mathbf{r};\omega) = \frac{\omega}{\pi} \sum_{\mu} \operatorname{Im} \left[\mathbf{\hat{n}}_{\alpha} \cdot \frac{\mathbf{E}_{\mu}(\mathbf{r}) \otimes \mathbf{E}_{\mu}(\mathbf{r})}{\tilde{\omega}_{\mu}(\tilde{\omega}_{\mu} - \omega)} \cdot \mathbf{\hat{n}}_{\alpha} \right]$$

See also: [C. Sauvan *et al.*, Phys. Rev. Lett. **110**, 237401 (2013)] [R.-C. Ge *et al.*, New J. Phys. **16**, 113048 (2014)]

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[P. T. Kristensen et al., Opt. Lett. 39, 6359 (2014)]

[J. R. de Lasson et al., J. Opt. Soc. Am. A 31, 2142 (2014)]

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$$p(\mathbf{r};\omega) = \frac{2\omega}{\pi c^2} \operatorname{Im} \left[\mathbf{\hat{n}}_{\alpha} \cdot \mathbf{G}(\mathbf{r},\mathbf{r};\omega) \cdot \mathbf{\hat{n}}_{\alpha} \right]$$

[L. Novotny and B. Hecht, "Principles of Nano-Optics" (2012)]

$$\rho(\mathbf{r};\omega) = \frac{\omega}{\pi} \sum_{\mu} \operatorname{Im} \left[\mathbf{\hat{n}}_{\alpha} \cdot \frac{\mathbf{E}_{\mu}(\mathbf{r}) \otimes \mathbf{E}_{\mu}(\mathbf{r})}{\tilde{\omega}_{\mu}(\tilde{\omega}_{\mu} - \omega)} \cdot \mathbf{\hat{n}}_{\alpha} \right]$$

See also: [C. Sauvan *et al.*, Phys. Rev. Lett. **110**, 237401 (2013)] [R.-C. Ge *et al.*, New J. Phys. **16**, 113048 (2014)]

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$$\rho^{y}(\mathbf{r}_{\mathrm{D}};\omega) = \frac{\omega}{\pi} \frac{1}{\epsilon(\mathbf{r}_{\mathrm{D}})} \mathrm{Im} \left[\frac{1}{\tilde{\omega}_{\mu}(\tilde{\omega}_{\mu}-\omega)} \frac{1}{a_{\mu}} \right]$$
$$a_{\mu} = \frac{1}{\epsilon(\mathbf{r}_{\mathrm{D}}) \left[\mathbf{E}_{\mu}(\mathbf{r}_{\mathrm{D}}) \cdot \hat{\mathbf{y}} \right]^{2}}$$

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$$\rho^{y}(\mathbf{r}_{\mathrm{D}};\omega) = \frac{\omega}{\pi} \frac{1}{\epsilon(\mathbf{r}_{\mathrm{D}})} \mathrm{Im} \left[\frac{1}{\tilde{\omega}_{\mu}(\tilde{\omega}_{\mu}-\omega)} \frac{1}{a_{\mu}} \right]$$
$$a_{\mu} = \frac{1}{\epsilon(\mathbf{r}_{\mathrm{D}}) \left[\mathbf{E}_{\mu}(\mathbf{r}_{\mathrm{D}}) \cdot \hat{\mathbf{y}} \right]^{2}}$$

$$F_{\rm P} \equiv \frac{\rho^y(\mathbf{r}_{\rm D};\omega_{\mu})}{\rho_{\rm Bulk}} \simeq \frac{1}{\pi^2} \left(\frac{\lambda_0}{n(\mathbf{r}_{\rm D})}\right)^2 \frac{Q_{\mu}}{A_{\mu}}$$

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$$\rho^{y}(\mathbf{r}_{\mathrm{D}};\omega) = \frac{\omega}{\pi} \frac{1}{\epsilon(\mathbf{r}_{\mathrm{D}})} \mathrm{Im} \left[\frac{1}{\tilde{\omega}_{\mu}(\tilde{\omega}_{\mu} - \omega)} \frac{1}{a_{\mu}} \right]$$
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[J. R. de Lasson et al., Opt. Lett. 40, 5790 (2015)]


















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00000000 2.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 $^{0}\gamma/z$ 1.5 0000000 Guided 0000000 00000000 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0.5 Frequency 0000000 Ó 0 0 0 0 0 0 0 Slow light -3 -2 -1 x/λ_0



PBG

Band edge

Zone folding

Normalized wavenumber

-requency

10 30 50

Group index

10























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 $(x_D, z_D) = (0, 0)$

 $-(x_D, z_D) = (0, -a/4)$

 ρ_G^{0}/ρ_{Bulk}

 $(x_D, z_D) = (0, -a/2)$

(a) = (100 nm (









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Nanophotonics Theory



Nanophotonics Theory



Nanophotonics



Nanophotonics Theory



Nanophotonics



EPFL-LPN



Nanophotonics Theory



Nanophotonics



EPFL-LPN



Niels



Jesper



Philip



September 3, 2007 \rightarrow December 18, 2015



September 3, 2007 \rightarrow December 18, 2015

...Thanks for your attention $\ddot{-}$

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